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## Fortran Codes for Computing the Space-Time Correlations of Turbulent Flow in a Channel

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<p>Computer codes have been developed which compute the space-time correlations in turbulent channel flow. The two-point spatial correlation function is computed by Fourier transforming the appropriate wavenumber spectrum. The wavenumber spectrum is computed by averaging over an appropriate number of ensembles and using geometric symmetries to improve the smoothness of the results. Space-time correlations are computed in a similar manner.</p> <p style="text-align: right;">(Continued on reverse side)</p>			
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## CONTENTS

1. INTRODUCTION .....	1
2. COMPUTATION OF THE COMPLETE TWO-POINT CORRELATION FUNCTION FOR TURBULENT CHANNEL FLOW FROM SPATIAL REALIZATIONS .....	2
3. CALCULATION OF THE SPACE-TIME CORRELATION FUNCTION FROM PLANAR DATA .....	9
CONCLUSIONS .....	13
ACKNOWLEDGEMENTS .....	14
REFERENCES .....	15
APPENDIX .....	17

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## FORTRAN PROGRAMS FOR COMPUTING THE SPACE-TIME CORRELATIONS OF TURBULENT FLOW IN A CHANNEL

### 1. Introduction

Recently, several groups have performed direct numerical simulations of turbulent flow in a channel at low Reynolds numbers. These have included the high resolution (fully resolved) calculation of Kim, Moin, and Moser (1987) and the low resolution calculations of Gilbert and Kleiser (1987). These simulations generate enormous quantities of data which have been used to investigate the validity of current theories of turbulence, to study mechanisms for the production of turbulence in wall-bounded flows, and to develop new turbulence models. However, before such detailed investigations can be undertaken, it is important to first check some basic statistics of the turbulence such as the mean velocity profile, turbulence intensities, and Reynolds stress profiles. Beyond these one-point statistics, it is also of considerable value to have a rapid means of computing the two-point statistics so that the length and time scales of the turbulence can be determined. In this report, we describe FORTRAN codes which we have developed which enable rapid evaluation of the two point correlations from a given dataset. Datasets were created in two forms. In the first, we have taken 'spatial snapshots' (realizations) of the flow. That is, at a given time in the calculation, we have stored the three components of the velocity at each collocation point designated by its coordinates  $x, y$ , and  $z$ , which are respectively the streamwise, spanwise, and wall-normal coordinates. In general, we have saved the realizations at time intervals on the order of or larger than the largest correlation time of the flow so that we may say that the dataset consists of a number of statistically independent snapshots. In the second type

of data, we have stored, at certain fixed locations above the wall, the three components of velocity in the horizontal planes at equally spaced time intervals. The time between saves for the planar data is generally significantly less than the correlation time of the flow. We will call such datasets ‘planar’ datasets. In Section 2 we describe codes which produce the two-point correlation function of turbulence from a ‘snapshot’ dataset and in Section 3 we describe the code which produces the space-time correlation function from a ‘planar’ dataset.

## 2. Computation of the Complete Two-Point Correlation Function for Turbulent Channel Flow from Spatial Realizations

### A. Mathematical Preliminaries and Definitions

Before describing the computer codes in detail, we briefly discuss certain mathematical definitions and the methods used to compute the two-point correlation function. (In the following discussions, we use the variables  $x_1, x_3$ , and  $x_2$  interchangeably with the notation  $x, y, z$ , and bold face notation indicates a vector quantity.) For any given flow field,  $\mathbf{u}(\mathbf{x}, t)$ , we obtain a set of statistically independent realizations by sampling the field at times separated by at least the characteristic correlation time of the flow. We define the Fourier coefficients of any such realization by:

$$\Phi(n, m, x_2) = \sum_{l=0}^{N_1-1} \sum_{k=0}^{N_3-1} \mathbf{u}(l, k, x_2) e^{\frac{2\pi n l i}{N_1}} e^{\frac{2\pi m k i}{N_3}},$$

where

$$i = \sqrt{-1}, \\ n = 0, \dots, \frac{N_1}{2} + 1, \\ m = -\frac{N_3}{2}, \dots, \frac{N_3}{2} - 1. \quad (1)$$

and

$$m = -\frac{N_3}{2}, \dots, \frac{N_3}{2} - 1.$$

The expression above defines collocation points  $(x_1)_l$  and  $(x_3)_k$ , and wavenumbers  $(k_1)_n$  and  $(k_3)_m$  as follows:

$$\begin{aligned}(x_1)_l &= l \times L_1/N_1, \\ (x_3)_k &= k \times L_3/N_3, \\ (k_1)_n &= n \times \frac{2\pi}{L_1}, \\ (k_3)_m &= m \times \frac{2\pi}{L_3},\end{aligned}\tag{2}$$

in which  $L_1$  and  $L_3$  are the domain lengths in the streamwise and spanwise directions respectively. In equation 1 time serves merely as a parameter which may be used to identify any given flow realization and is therefore suppressed. We now follow Sirovich (1987), who suggests that the geometric symmetries of a flow should be exploited to increase the effective number of realizations that are used in the averaging process. For the case of channel flow, the equations of motion are invariant with respect to wall normal reflection, spanwise reflection, and 180 degree rotation about the streamwise axis. Applying these symmetries to the spectral representation  $\Phi^0$ , where we use the superscript 0 to refer to the representation before reflections and rotations have been applied, yields three additional representations of the flow which are given by:

$$\begin{aligned}\Phi^1 &= \{\Phi_1^0(n, m, -x_2), -\Phi_2^0(n, m, -x_2), \Phi_3^0(n, m, -x_2)\} \\ \Phi^2 &= \{\Phi_1^0(n, -m, x_2), \Phi_2^0(n, -m, x_2), -\Phi_3^0(n, -m, x_2)\}, \\ \Phi^3 &= \{\Phi_1^0(n, -m, -x_2), -\Phi_2^0(n, -m, -x_2), -\Phi_3^0(n, -m, -x_2)\}.\end{aligned}\tag{3}$$

In these expressions, the superscripts 1, 2, and 3 refer to vertical reflection, spanwise reflection, and streamwise rotation symmetries respectively. The subscripts continue to represent the coordinate directions as previously defined. The (complex) spectrum  $\Psi_{\alpha\beta}$  can now be computed from:

$$\Psi_{\alpha\beta}(n, m, x_2, x_2') = \langle \overline{\frac{1}{4} \sum_{p=0}^3 \Phi_\alpha^p(n, m, x_2) \Phi_\beta^p(n, m, x_2')} \rangle, \quad (4)$$

where

$$\alpha, \beta = 1, 2, 3.$$

The overbar designates complex conjugate, and the brackets represent an average over all the available realizations. We note that in addition to the increase in the effective size of the data base by a factor of four, the symmetries have the additional advantage of reducing the size of  $\Psi_{\alpha\beta}$  also by a factor of four. A factor of two comes from the spanwise symmetry which insures that  $\Psi_{\alpha\beta}$  is unique for  $n > 0$  and  $m > 0$ . A second factor of two comes from the vertical reflection symmetry which gives unique values of the spectrum for  $-1 < x_2 < 0$  and  $-1 < x_2' < 1$ . The two point correlation function,  $R_{\alpha\beta}$ , can now be obtained from:

$$R_{\alpha\beta}(l, k, x_2, x_2') = K \sum_{n=0}^{\frac{N_1}{2}-1} \sum_{m=-\frac{N_3}{2}}^{\frac{N_3}{2}-1} \Psi_{\alpha\beta}(n, m, x_2, x_2') e^{\frac{2\pi nl i}{N_1}} e^{\frac{2\pi mk i}{N_3}}, \quad (5)$$

where

$$K = \frac{1}{\sqrt{R_{\alpha\alpha}(0, 0, x_2) R_{\beta\beta}(0, 0, x_2')}}.$$

We note that care must be taken in the calculation of  $R_{\alpha\beta}$  since  $\Psi_{\alpha\beta}$  is computed only for  $n, m \geq 0$ . That is, the inherent symmetries in  $\Psi_{\alpha\beta}$  must be taken into account when equation 5 is used. In particular, the following symmetries must hold:

$$\begin{aligned}\Psi_{11}(n, -m, x_2, x'_2) &= \Psi_{11}(n, m, x_2, x'_2), \\ \Psi_{22}(n, -m, x_2, x'_2) &= \Psi_{22}(n, m, x_2, x'_2), \\ \Psi_{33}(n, -m, x_2, x'_2) &= \Psi_{33}(n, m, x_2, x'_2), \\ \Psi_{12}(n, -m, x_2, x'_2) &= \Psi_{12}(n, m, x_2, x'_2), \\ \Psi_{13}(n, -m, x_2, x'_2) &= -\Psi_{13}(n, m, x_2, x'_2), \\ \Psi_{23}(n, -m, x_2, x'_2) &= -\Psi_{23}(n, m, x_2, x'_2).\end{aligned}\tag{6}$$

A detailed check on the computation of  $\Psi_{\alpha\beta}$  was made by computing the energies (mean square values) of each velocity component by summing  $\Psi_{\alpha\beta}$  over  $m$  and  $n$ . The comparison of these values with the known energies was exact and will be described in a future report.

### B. Description of the codes KLEXP and KLANA

The computation of the correlation function  $R_{\alpha\beta}$  defined by equation 5 is computed by a two-stage process. First, the spectral function  $\Psi_{\alpha\beta}$  defined by equation 4 is computed and saved. This computation is performed using the code KLEXP. Then the spectral function is used to calculate  $R_{\alpha\beta}$  using the code KLANA. Originally, the interest was primarily to perform an eigenfunction decomposition (Lumley(1970)) of turbulent flow in a channel. To perform such a decomposition, the code KLEXP was written to determine  $\Psi_{\alpha\beta}$  which was of greatest interest for our purposes. Later it was decided that the real space function  $R_{\alpha\beta}$  would also be of interest so that the code

KLANA was written for the purpose of computing it. All codes discussed in this report are listed in the Appendix.

### B.1 Description of the code KLEXP

The code KLEXP listed in the Appendix was written to compute  $\Psi_{\alpha\beta}$  from a data set which consists of 30 realizations of channel flow turbulence. In this particular case, there were 7 data files containing 4 realizations and one containing 2. On each file, the three components of vorticity were also stored but were skipped each time the data was accessed. After each realization of the velocity field is read into memory, certain baseline statistics such as the mean and variance of each velocity component are computed. These values, which are calculated from the raw (real space) data, can later be compared to the values obtained from the spectral function  $\Psi_{\alpha\beta}$  as a check on the calculation. Next, the horizontal transform of the three components of velocity is performed using the subroutine YXFOUR. This calculation is represented mathematically by equation 1. The six components of  $\Psi_{\alpha\beta}$  are then calculated from the raw spectral representation,  $\Phi_\alpha$ , using the subroutine MATRIX. In this routine, for each spectral realization of the flow, the product  $\overline{\Phi_\alpha}\Phi_\beta$  is computed and averaging is performed over the symmetry groups given in equation 3. We note that the parameter ‘is’ is used to distinguish the four possible ways (which will be described below) in which the symmetries need to be applied.

We now describe in detail how the six components of  $\Psi_{\alpha\beta}$  are computed. In computing  $\Psi_{11}$  we first compute  $r_{11}^0$  and  $r_{11}^1$  which are given by:

$$r_{11}^0(n, m, x_2, x_2') = \overline{\Phi_1}(n, m, x_2)\Phi_1(n, m, x_2'), \quad (7)$$

and

$$r_{11}^1(n, m, x_2, x_2') = \overline{\Phi_1}(n, m, -x_2)\Phi_1(n, m, -x_2').$$

We then apply the spanwise reflection symmetry to  $r_{11}^0$  and  $r_{11}^1$  using the subroutine SYM1. This yields  $r_{11}^2$  and  $r_{11}^3$  which are given by:

$$r_{11}^2(n, m, x_2, x_2') = r_{11}^0(n, m, x_2, x_2') + r_{11}^0(n, -m, x_2, x_2'),$$

and

$$r_{11}^3(n, m, x_2, x_2') = r_{11}^1(n, m, x_2, x_2') + r_{11}^1(n, -m, x_2, x_2'). \quad (8)$$

Finally, the function  $\Psi_{11}$  is computed from:

$$\Psi_{11}(n, m, x_2, x_2') = r_{11}^2(n, m, x_2, x_2') + r_{11}^3(n, m, x_2, x_2'). \quad (9)$$

This procedure incorporates all three group symmetries given in equation 3. The same procedure is used to compute  $\Psi_{22}$  and  $\Psi_{33}$ .

To obtain the other three components of the correlation matrix, minor modifications of the procedure outlined above are needed. In the computation of  $\Psi_{12}$  we must take into account the change in sign of the wall-normal velocity. This is expressed by :

$$\Psi_{12}(n, m, x_2, x_2') = r_{12}^2(n, m, x_2, x_2') - r_{12}^3(n, m, x_2, x_2'). \quad (10)$$

In the computation of  $\Psi_{13}$  we must modify the relations in equation 8 since there is a sign change in the spanwise velocity component when the spanwise reflection symmetry is used. The resulting expressions for  $r_{13}^2$  and  $r_{13}^3$  are given by:

$$r_{13}^2(n, m, x_2, x_2') = r_{13}^0(n, m, x_2, x_2') - r_{13}^0(n, -m, x_2, x_2'),$$

and

$$r_{13}^3(n, m, x_2, x_2') = r_{13}^1(n, m, x_2, x_2') - r_{13}^1(n, -m, x_2, x_2'). \quad (11)$$

In calculating  $\Psi_{23}$  we must take into account both the sign changes in the wall-normal velocity component and the spanwise velocity component. This results in the following relations:

$$\begin{aligned} r_{23}^2(n, m, x_2, x_2') &= r_{23}^0(n, m, x_2, x_2') - r_{23}^0(n, -m, x_2, x_2'), \\ r_{23}^3(n, m, x_2, x_2') &= r_{23}^1(n, m, x_2, x_2') - r_{23}^1(n, -m, x_2, x_2'), \end{aligned} \quad (12)$$

and

$$\Psi_{23}(n, m, x_2, x_2') = r_{23}^2(n, m, x_2, x_2') - r_{23}^3(n, m, x_2, x_2'). \quad (13)$$

The final result,  $\Psi_{\alpha\beta}$ , is then scaled appropriately and diagnostics are applied to check the answer. A good check is to determine whether the energy in  $\Psi_{\alpha\beta}$ , obtained by integrating over all wavenumbers, is equal to the energies computed from the raw (real space) fields. The final result is then saved on disc. No unnecessary information has been computed or stored in this calculation and, as noted previously, the application of the symmetries reduces the storage requirements by a factor of four.

## B.2 Description of the code KLANA

Code KLANA uses the spectral results  $\Psi_{\alpha\beta}$  produced by KLEXP as described in equation 4 to compute the correlations  $R_{\alpha\beta}$  described in equation 5. Before the main calculation, all six spectral functions are read from disc. The main loop (do 8000) computes the correlation functions one at a time and writes each result to disc successively. Since the calculation of each of the six correlations is identical except for the application of the symmetries described in equation 6, a description of the calculation of  $R_{11}$  will serve as a description for the other five. For each pair of indices 'iz' and 'izp' which designate wall-normal distances, a two dimensional array 'phi' is extracted from the four dimensional spectral function  $\Phi_{11}$ . Since the resultant correlation must

be real, the known symmetry properties of  $\Phi_{11}$  are applied to the array 'phi' in subroutine SETPHI to generate the array 'phisym'. The two dimensional Fourier transform is then performed in subroutine RIJ to complete the calculation. Note that the parameter 'ij' in the call to RIJ determines which of the symmetries given in equation 6 will be applied. The final result is placed into the larger array 'rrij', which is then normalized using the proper  $K$  defined in equation 5. These correlation functions and the normalization factors which have been used are written to a disk file for later plotting.

### 3. Calculation of the Space-Time Correlation Function from Planar Data

#### A. Mathematical Definitions

Before describing the details of the FORTRAN code that has been written to compute the space-time correlation function, we first describe the method of calculation. Given the planar realization  $\mathbf{u}(x_1, x_3, t)$ , we define the following spectral representations:

$$\begin{aligned}\gamma_\alpha(n, m, t) &= \sum_{l=0}^{N_1-1} \sum_{k=0}^{N_3-1} u_\alpha(l, k, t) e^{\frac{2\pi nl_i}{N_1}} e^{\frac{2\pi mk_i}{N_3}}, \\ \theta_\alpha(n, j, x_3) &= \sum_{l=0}^{N_1-1} \sum_{k=0}^{N_t-1} u_\alpha(l, k, x_3) e^{\frac{2\pi nl_i}{N_1}} e^{\frac{2\pi kj_i}{N_t}}, \\ \pi_\alpha(m, j, x_1) &= \sum_{l=0}^{N_3-1} \sum_{k=0}^{N_t-1} u_\alpha(l, k, x_1) e^{\frac{2\pi ml_i}{N_3}} e^{\frac{2\pi kj_i}{N_t}},\end{aligned}\quad (14)$$

and

$$\alpha = 1, 2, 3.$$

We note that  $x_2$  is suppressed in equation 14 since it is assumed to be fixed in these calculations. We also define the (discretized) time and frequency as follows:

$$t_n = \frac{T}{N_t} n,$$

$$\omega_n = \frac{2\pi}{T}n \quad (15)$$

and

$$n = 0, 1, \dots, N_t - 1.$$

$T$  is the total time in any given realization. In these calculations we apply the spanwise symmetry to  $\gamma$  and  $\pi$  in the following way:

$$\gamma^1 = \{\gamma_1^0(n, -m, t), \gamma_2^0(n, -m, t), -\gamma_3^0(n, -m, t)\},$$

and

$$\pi^1 = \{\pi_1^0(-m, j, t), \pi_2^0(-m, j, t), -\pi_3^0(-m, j, t)\}. \quad (16)$$

where the subscripts and superscripts have been defined in Section 2. We can then compute the (complex) spectra defined by:

$$\begin{aligned} \Gamma_{\alpha\beta}(n, m) &= \left\langle \frac{1}{N_t} \sum_{k=0}^{N_t-1} \frac{1}{2} \sum_{p=0}^1 \overline{\gamma_\alpha^p(n, m, t_k)} \gamma_\beta^p(n, m, t_k) \right\rangle, \\ \Theta_{\alpha\beta}(n, j) &= \left\langle \frac{1}{N_3} \sum_{k=0}^{N_3-1} \overline{\theta_\alpha^p(n, m, (x_3)_k)} \theta_\beta^p(n, m, (x_3)_k) \right\rangle, \\ \Pi_{\alpha\beta}(m, j) &= \left\langle \frac{1}{N_1} \sum_{k=0}^{N_1-1} \frac{1}{2} \sum_{p=0}^1 \overline{\pi_\alpha^p(m, j, (x_1)_k)} \pi_\beta^p(m, j, (x_1)_k) \right\rangle, \end{aligned} \quad (17)$$

and

$$\alpha, \beta = 1, 2, 3.$$

We note again that no symmetry has been applied in the calculation of  $\Theta_{\alpha\beta}$  as indicated in equation 17. In summary, we compute the appropriate spectral functions by first averaging over the appropriate symmetries, then averaging over a coordinate, and finally averaging over the available realizations.

We then define, as before, the two point correlation functions by:

$$\begin{aligned}
 R^a_{\alpha\beta}(l, k) &= K_a \sum_{n=0}^{\frac{N_1}{2}-1} \sum_{m=-\frac{N_3}{2}}^{\frac{N_3}{2}-1} \Gamma_{\alpha\beta}(n, m) e^{\frac{2\pi n li}{N_1}} e^{\frac{2\pi m ki}{N_3}}, \\
 R^b_{\alpha\beta}(l, k) &= K_b \sum_{n=0}^{\frac{N_1}{2}-1} \sum_{j=0}^{N_t-1} \Theta_{\alpha\beta}(n, j) e^{\frac{2\pi n li}{N_1}} e^{\frac{2\pi j ki}{N_t}}, \\
 R^c_{\alpha\beta}(l, k) &= K_c \sum_{m=-\frac{N_3}{2}}^{\frac{N_3}{2}-1} \sum_{j=0}^{N_t-1} \Pi_{\alpha\beta}(m, j) e^{\frac{2\pi m li}{N_3}} e^{\frac{2\pi j ki}{N_t}},
 \end{aligned} \tag{18}$$

and

$$\begin{aligned}
 K_a &= \frac{1}{\sqrt{R^a_{\alpha\alpha}(0, 0)R^a_{\beta\beta}(0, 0)}}, \\
 K_b &= \frac{1}{\sqrt{R^b_{\alpha\alpha}(0, 0)R^b_{\beta\beta}(0, 0)}}, \\
 K_c &= \frac{1}{\sqrt{R^c_{\alpha\alpha}(0, 0)R^c_{\beta\beta}(0, 0)}}.
 \end{aligned} \tag{19}$$

## B. Description of code WAVSPEC

The Fortran code WAVSPEC permits the analysis of data of the 'planar' type as described in the Introduction. This data consists of the three components of velocity  $u_1$ ,  $u_2$ , and  $u_3$ , the streamwise, spanwise, and normal components of the vector  $\mathbf{u}(x_1, x_3, t)$ . The code listed in the Appendix was written to compute correlations from a dataset consisting of seven files in which horizontal planes of the three components of velocity and the vorticity are stored. Each file contains 512 time steps. After each realization (i.e. a file consisting of 512 time steps) is read from disc using the subroutine READ, the computations performed in equations 14-17 are performed by subroutine HFENS.

(Array 'tmp' in subroutine READ is a dummy used to skip over the values of vorticity contained in the dataset.)

The input to subroutine HFENS is the raw planar data, and the output consists of the updated (ensemble averaged) spectra given by  $\Gamma$ ,  $\Theta$ , and  $\Pi$  defined in equation 17. We describe only the calculation of  $\Gamma$  in HFENS, since the calculation of  $\Theta$ , and  $\Pi$  is virtually the same. The first loop in HFENS (do 100) computes the average over time as indicated in equation 17. The second loop (do 70) extracts the two dimensional arrays (the y-x array) from the three dimensional data (y-x-t arrays). Each two dimensional array is then brought into wavenumber ( $k_3 - k_1$ ) space using the subroutine HFFT1. The subroutine ENSAV1 then computes the square of the magnitude of the Fourier coefficients and applies the spanwise reflection symmetry. This completes the calculation of  $\Gamma$ . As indicated in equation 17, no symmetry is applied in the the calculation of  $\Theta$ .

As a check on the calculation, we observe that if we sum  $\Gamma$ ,  $\Theta$ , and  $\Pi$  over wavenumber and frequency, the results must be identical. This sum is the mean square value (the energy) of the velocity component. This serves as a check on the calculations and therefore, following the action of HFENS, the total energy for each of the three spectra are computed using the routines ENRG1, ENRG2, and ENRG3. The results of these calculations, which will be given in a subsequent report, show that the energies were equal to within machine roundoff error. After the energies have been computed, they are used to normalize the spectral results.

The spectra  $\Gamma$ ,  $\Theta$ , and  $\Pi$  are then used to compute the space-time or space-space correlations given in equation 18. This is accomplished by first calling SETPH1, SETPH2, and SETPH3, each of which arranges the spectral data in an appropriate form for subsequent Fourier transformation. The Fourier transforms are performed

by the routines RIJ1, RIJ2, and RIJ3, the correlations are normalized as indicated in equation 18, and the results are saved on disc.

### Conclusions

We have described computer codes that have been written to compute the relevant correlation functions and their corresponding spectral representations from stochastic velocity data which has been obtained from direct simulations of channel flow turbulence. The codes we have written can be used to analyze data that is in one of two forms : (1) spatial 'snapshots' of the flow , and (2) 'planar' datasets. In the first case we compute the relevant correlations in a two step process. First, we compute the four-dimensional spectrum given in equation 4 in which we average over all available realizations and also incorporate the relevant symmetries. This intermediate result is of particular importance since it forms the basis for a Karhunen-Loeve or eigenfunction expansion to which we referred in Section 2. This calculation is performed by the code KLEXP. These results are then used to compute the (real-space) correlation function given in equation 5 using the code KLANA. For the second case, using data of the 'planar' form, the relevant spectra given in equation 17 and correlations given in equation 18 are computed by implementing the code WAVSPEC.

As we mentioned in the Introduction, along with the theoretical interest that exists in computing these multidimensional statistics, these statistics also serve as diagnostic tools which can be used to check the very large data bases that are generated from direct numerical simulations. A description and some preliminary interpretations of the results of these calculations will be given in a subsequent report.

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## Appendix

Here we list the codes KLEXP, KLANA, and WAVSPEC described in the text. It is hoped that the text description together with the annotations within each code will enable a user to make modifications for his particular application. Note that the fft routines CFFT2, RCFFT2, and CRFFT2 are library subroutines which are generally available on the Cray-XMP.



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        read(lu) (((om1(k,j,i),k=1,nz1),j=1,1),i=1,1)
CCCCCCCCCCCCCCC COMPUTE MEAN VALUES CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
CCCC THIS IS BEING DONE MERELY FOR DIAGNOSTIC PURPOSES. THAT IS,
CC WE ARE COMPUTING THE MEAN AND RMS FOR EACH OF THE THREE FIELDS.
        call mean(u1,u1m)
        call mean(u2,u2m)
        call mean(u3,u3m)
        call zms(u1,u1,u1rms)
        call zms(u2,u2,u2rms)
        call zms(u3,u3,u3rms)
        call zms(u1,u2,u12rms)
        call zms(u1,u3,u13rms)
        call zms(u2,u3,u23rms)
CCCCCCCCCCCCCCC GO TO FOURIER SPACE CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
CCCC TRANSFORM EACH OF THE FIELDS FROM U(Z,Y,X) TO U(Z,M,N)
CCCC BUT FIRST LETS INITIALIZE THE FFTS.
        call expo
        call yxfour(u1)
        call yxfour(u2)
        call yxfour(u3)
CCCCCCCCCCCCCCC COMPUTE CORRELATION MATRICRES CCCCCCCCCCCCCCCCC
C COMPUTE EACH OF THE SIX INDEPENDENT MEMBERS OF THE CORRELATION
C MATRIX
        call matrix(u1,u1,ulu1,1)
        call matrix(u2,u2,u2u2,1)
        call matrix(u3,u3,u3u3,1)
        call matrix(u1,u2,ulu2,2)
        call matrix(u1,u3,ulu3,3)
        call matrix(u2,u3,u2u3,4)
1      continue
CCCCCCCCCCCCCCC NOW SCALE RESULTS CCCCCCCCCCCCCCCCCCCCCCCCCCCCC
        con = 1./(float(30)*fac)
        con1 = 1./(float(nens)*fac1)
        call scale(ulu1,con1)
        call scale(u2u2,con1)
        call scale(u3u3,con1)
        call scale(ulu2,con1)
        call scale(ulu3,con1)
        call scale(u2u3,con1)
        call mult(u1m,con)
        call mult(u2m,con)
        call mult(u3m,con)
        call mult(u1rms,con)
        call mult(u2rms,con)
        call mult(u3rms,con)
        call mult(u12rms,con)
        call mult(u13rms,con)
        call mult(u23rms,con)
CCCCCCCCCCCCCCC COMPUTE RMS VALUES CCCCCCCCCCCCCCCCCCCCCCCCCCCCC
        do 2 k = 1,nz1
            u1rms(k) = sqrt(u1rms(k) - u1m(k)**2)
            u2rms(k) = sqrt(u2rms(k) - u2m(k)**2)
            u3rms(k) = sqrt(u3rms(k) - u3m(k)**2)
            u12rms(k) = (u12rms(k)- u1m(k)*u2m(k))
            u13rms(k) = (u13rms(k)- u1m(k)*u3m(k))
            u23rms(k) = (u23rms(k)- u2m(k)*u3m(k))
2      continue
            u12rms(1) = 0.0
            u12rms(nz1) = 0.0
            u13rms(1) = 0.0

```

```

        u13rms(nz1) = 0.0
        u23rms(1) = 0.0
        u23rms(nz1) = 0.0
    do 3 k =2,nz1-1
        u12rms(k) = u12rms(k)/(ulrms(k)*u2rms(k))
        u13rms(k) = u13rms(k)/(ulrms(k)*u3rms(k))
        u23rms(k) = u23rms(k)/(u2rms(k)*u3rms(k))
3         continue
CCCCCCCCCCCCCCCC COMPUTE RMS VALUES FROM SPECTRAL RESULTS CCCCCCCCCCCCCCCCC
C      HERE WE INTEGRATE THE CORRELATION FUNCTIONS (AFTER HAVING SCALED
C      THEM PROPERLY) TO SEE IF WE PERFORMED THE CALCULATION PROPERLY.
C      THAT IS, DOES THE RMS WE GET FROM THE SPECTRA EQUAL THE TRUE RMS
C      FOR EACH FIELD?
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
    call intg(ul1,s1)
    call intg(u2u2,s2)
    call intg(u3u3,s3)
    call intg(ulu2,s12)
    do 4 k =1,nzd21
        s1(k) = sqrt(2.0 * s1(k))
        s2(k) = sqrt(2.0 * s2(k))
4         s3(k) = sqrt(2.0 * s3(k))
        s12(1) = 0.0
        do 5 k = 2,nzd21
            s12(k) = (2.* s12(k))/(s1(k)*s2(k))
5         continue
CCCCCCCCCCCCCCCCCCCC WRITE RESULTS TO DISC CCCCCCCCCCCCCCCCCCCCCCCCCCCCC
C      FIRST DUMP THE CORRELATION MATRICES
C      NOTE WE ARE NOW DUMPING FROM IZ =1,NZD21
C      AND IZP = 1,NZ1. THAT IS ONE HALF OF EACH MATRIX.
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
    do 10 izp =1,nz1
        write(50) (((ul1(j,i,iz,izp),j=1,23),i=1,14),iz=1,nzd21)
        write(51) (((u2u2(j,i,iz,izp),j=1,23),i=1,14),iz=1,nzd21)
        write(52) (((u3u3(j,i,iz,izp),j=1,23),i=1,14),iz=1,nzd21)
        write(53) (((ulu2(j,i,iz,izp),j=1,23),i=1,14),iz=1,nzd21)
        write(54) (((ulu3(j,i,iz,izp),j=1,23),i=1,14),iz=1,nzd21)
        write(55) (((u2u3(j,i,iz,izp),j=1,23),i=1,14),iz=1,nzd21)
10       continue
C      THEN DUMP THE MEAN VALUES AND THE RMS VALUES OF EACH FIELD
        write(60,100)(ulm(k),k=1,nz1)
        write(61,100)(u2m(k),k=1,nz1)
        write(62,100)(u3m(k),k=1,nz1)
        write(63,100)(ulrms(k),k=1,nz1)
        write(64,100)(u2rms(k),k=1,nz1)
        write(65,100)(u3rms(k),k=1,nz1)
        write(66,100)(u12rms(k),k=1,nz1)
        write(67,100)(u13rms(k),k=1,nz1)
        write(68,100)(u23rms(k),k=1,nz1)
CCC     THEN DUMP THE MEAN SQUARE VALUES COMPUTED FROM THE SPECTRA.
C      NOTE THAT SINCE THE SYMMETRIES WERE USED TO IN COMPUTING
C      THESE SPECTRA THE DEGENERACIES MUST BE INCLUDED IN THE
C      COMPUTATION OF THE ENERGIES. THIS HAS NOT BEEN DONE HERE
C      SO THAT THE ENERGY VALUES WE GET HERE WILL NOT AGREE
C      WITH THE VALUES WE GET FROM THE RAW DATA. THESE DEGENERACIES
C      CAN BE ACCOUNTED FOR AND THE RESULTS WHEN WE DO INCLUDE THEM
C      PROPERLY DO APEAR TO BE CORRECT.
        write(70,100)(s1(k),k=1,nzd21)
        write(71,100)(s2(k),k=1,nzd21)
        write(72,100)(s3(k),k=1,nzd21)

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        write(73,100)(s12(k),k=1,nzd21)
cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
100  format(2x,4e12.5)
      stop
      end
      subroutine matrix(a,b,c,is)
      parameter(nz=32,ny=64,nx=16)
      parameter(nz1=nz+1,nzml=nz-1,nzd21=nz/2+1,nzd2=nz/2)
      parameter(nyd2=ny/2,nyd21=ny/2+1)
      parameter(nxp2=nx+2,nxd2=nx/2,nxd21=nx/2+1)
      dimension a(nz1,ny,nxp2),b(nz1,ny,nxp2),c(nyd2,nxp2,nzd21,nz1)
      dimension r1(ny,nxp2),r2(ny,nxp2)

C      IN THIS SUBROUTINE WE COMPUTE THE CORRELATION (C) FROM THE
C      FIELDS A AND B. WE UTILIZE ALL THE AVAILABLE SYMMETRIES (4) FOR
C      THE CHANNEL FLOW PROBLEM. FOR EACH PAIR OF WALL NORMAL INDICES
C      IZ AND IZP WE COMPUTE THEIR COUNTERPARTS IZ1 AND IZ2. THE INDEX
C      IS DETERMINES THE SIGNS TO BE USED WHEN THE SYMMETRIES ARE APPLIED.
C      THE SUBROUTINES SYM1 AND SYM2 ARE USED TO APPLY THE REFLECTIONAL
C      SYMMETRIES. PLEASE NOTE!!! IZP GOES FROM 1 TO NZ1.
      do 1 iz = 1,nzd21
          iz1 = nz1 - iz + 1
      do 1 izp = 1,nz1
          iz2 = nz1 - izp + 1
do 2   j=1,ny
do 2   i=1,nxd21
      ir = 2*i -1
      im = 2*i
      wlr = a(iz,j,ir)
      wlim = a(iz,j,im)
      w2r = b(izp,j,ir)
      w2im = b(izp,j,im)
      w1pr = a(iz1,j,ir)
      w1pim = a(iz1,j,im)
      w2pr = b(iz2,j,ir)
      w2pim = b(iz2,j,im)
      r1(j,ir) = wlr*w2r + wlim*w2im
      r1(j,im) = wlr*w2im - wlim*w2r
      r2(j,ir) = w1pr*w2pr + w1pim*w2pim
      r2(j,im) = w1pr*w2pim - w1pim*w2pr
2      continue
CCCCCCCCCCCCCCCCCCCCCCCC      INVOKE SYMMETRIES CCCCCCCCCCCCCCCCCCCCCCCCC
      if ( is .eq. 1 ) then
          call sym1(r1,ny,nxp2)
          call sym1(r2,ny,nxp2)
      do 3 i=1,nxp2
      do 3 j=1,nyd2
          r1(j,i) = r1(j,i) + r2(j,i)
          endif
      if ( is .eq. 2 ) then
          call sym1(r1,ny,nxp2)
          call sym1(r2,ny,nxp2)
      do 4 i=1,nxp2
      do 4 j=1,nyd2
          r1(j,i) = r1(j,i) - r2(j,i)
          endif
      if ( is .eq. 3 ) then
          call sym2(r1,ny,nxp2)
          call sym2(r2,ny,nxp2)
      do 5 i=1,nxp2
      do 5 j=1,nyd2

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```

5           r1(j,i) = r1(j,i) + r2(j,i)
           endif
           if ( is .eq. 4 ) then
               call sym2(r1,ny,nxp2)
               call sym2(r2,ny,nxp2)
           do 6 i=1,nxp2
           do 6 j=1,nyd2
6           r1(j,i) = r1(j,i) - r2(j,i)
           endif
CCCCCCCCCCCCC   CALCULATE ENSEMBLE AVERAGE CCCCCCCCCCCCCCCCCCCCCCCCC
           do 7 i = 1,nxp2
               do 7 j = 1,nyd2
                   c(j,i,iz,izp) = c(j,i,iz,izp) + r1(j,i)
7           continue
CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC
1           continue
           return
           end
           subroutine setup
           parameter(ny=64,nx=16)
           parameter(nxp2=nx+2,nxd21=nx/2+1)
           parameter(nyt2=2*ny,nyd21=ny/2+1,nyd22=ny/2+2)
           common/param/fac,fac1
CCCCCCC ----- some parameters -----
           fac = float(nx*ny)
           fac1 = ((fac**2)*2)
           return
           end
           subroutine hfft(a,is)
           parameter(ny=64,nx=16)
           parameter(nxp2=nx+2,nxd21=nx/2+1)
           parameter(nyt2=2*ny,nyd21=ny/2+1,nyd22=ny/2+2)
           parameter(nx3p4=3*nx+4,ny5=5*ny,nxny=(nx+2)*ny)
           dimension a(ny,nxp2)
           common/index/ indx(nxp2),indy(nyt2)
           common/exp/ exrc(nx3p4),excr(nx3p4),eycc(ny5)
           common/work/ vecx(nxp2),vecy(nyt2),vecyl(nyt2)
               if (is.eq. -1) go to 500
           do 1 j=1,ny
           do 2 i=1,nxp2
               indx(i)=(i-1)*ny+j
2           continue
           call gather(nx,vecx,a,indx)
           call rcfft2(0,1,nx,vecx,exrc,vecx)
           call scatter(nxp2,a,indx,vecx)
1           continue
           do 3 i=1,nxd21
           do 4 j=1,ny
               j1=2*(j-1)+1
               indy(j1)=j+2*(i-1)*ny
               indy(j1+1)=indy(j1)+ny
4           continue
           call gather(nyt2,vecy,a,indy)
           call cfft2(0,1,ny,vecy,eycc,vecy)
           call scatter(nyt2,a,indy,vecy)
3           continue
           return
500           continue
           do 5 i=1,nxd21
           do 6 j=1,ny

```

```

        j1=2*(j-1)+1
        indy(j1)=j+2*(i-1)*ny
        indy(j1+1)=indy(j1)+ny
6      continue
        call gather(nyt2,vecy,a,indy)
        call cfft2(0,-1,ny,vecy,eycc,vecy)
        call scatter(nyt2,a,indy,vecy)
5      continue
        do 7 j=1,ny
        do 8 i=1,nxp2
            indx(i)=(i-1)*ny+j
8      continue
            call gather(nxp2,vecx,a,indx)
            call crfft2(0,-1,nx,vecx,excr,vecx)
            call scatter(nx,a,indx,vecx)
7      continue
        do 10 i=1,nx
        do 10 j=1,ny
            a(j,i)=a(j,i)/(2.*float(nx*ny))
10     continue
        return
        end
subroutine sym1(a,ny,nx)
dimension a(ny,nx)
nyd2=ny/2
nyd21 =ny/2+1
    do 1 i=1,nx
    do 1 j=1,nyd2
        if ( j .eq. 1) then
            a(j,i) = 2.0*a(j,i)
        else
            a(j,i) = (a(j,i) + a(ny-j+2,i))
        endif
1      continue
    return
    end
subroutine sym2(a,ny,nx)
CCCCCCC NOTE: BECAUSE OF THE SYMMETRY IN THIS CASE THE ZERO SPANWISE
CCCCCCC MODES MUST BE ZERO
dimension a(ny,nx)
nyd2=ny/2
nyd21 =ny/2+1
    do 1 i=1,nx
    do 1 j=1,nyd2
        if ( j .eq. 1) then
            a(j,i) = 0.0
        else
            a(j,i) = (a(j,i) - a(ny-j+2,i))
        endif
1      continue
    return
    end
subroutine scale(a,con)
parameter(nz=32,ny=64,nx=16)
parameter(nz1=nz+1,nzm1=nz-1,nzd21=nz/2+1,nzd2=nz/2)
parameter(nyd2=ny/2,nyd21=ny/2+1)
parameter(nxp2=nx+2,nxd2=nx/2,nxd21=nx/2+1)
dimension a(nyd2,nxp2,nzd21,nz1)
    do 1 izp = 1,nz1
    do 1 iz = 1,nzd21

```

```

        a(1,1,iz,izp) = 0.0
do 1 i=1,nxp2
do 1 j=1,nyd2
1      a(j,i,iz,izp) = a(j,i,iz,izp)*con
      return
      end
      subroutine mult(a,con)
parameter(nz=32,ny=64,nx=16)
parameter(nz1=nz+1,nzml=nz-1,nzd21=nz/2+1,nzd2=nz/2)
parameter(nyd2=ny/2,nyd21=ny/2+1)
parameter(nxp2=nx+2,nxd2=nx/2,nxd21=nx/2+1)
dimension a(nz1)
do 1 iz = 1,nz1
1      a(iz) = a(iz)*con
      return
      end
      subroutine yxfour(a)
parameter(nz=32,ny=64,nx=16)
parameter(nz1=nz+1,nzml=nz-1,nzd21=nz/2+1,nzd2=nz/2)
parameter(nyd2=ny/2,nyd21=ny/2+1)
parameter(nxp2=nx+2,nxd2=nx/2,nxd21=nx/2+1)
dimension a(nz1,ny,nxp2),work(ny,nxp2)
      do 1 k =1,nz1
          do 2 j =1,ny
              do 2 i =1,nx
                  work(j,i) = a(k,j,i)
2      call hfft(work,1)
          do 3 j =1,ny
              do 3 i =1,nxp2
                  a(k,j,i) = work(j,i)
3      continue
1      return
      end
      subroutine mean(a,s)
parameter(nz=32,ny=64,nx=16)
parameter(nz1=nz+1,nzml=nz-1,nzd21=nz/2+1,nzd2=nz/2)
parameter(nyd2=ny/2,nyd21=ny/2+1)
parameter(nxp2=nx+2,nxd2=nx/2,nxd21=nx/2+1)
dimension a(nz1,ny,nxp2),s(nz1)
      do 2 i =1,nx
          do 2 j =1,ny
              do 2 k =1,nz1
                  s(k) = a(k,j,i) + s(k)
2      return
      end
      subroutine zms(a,b,s)
parameter(nz=32,ny=64,nx=16)
parameter(nz1=nz+1,nzml=nz-1,nzd21=nz/2+1,nzd2=nz/2)
parameter(nyd2=ny/2,nyd21=ny/2+1)
parameter(nxp2=nx+2,nxd2=nx/2,nxd21=nx/2+1)
dimension a(nz1,ny,nxp2),b(nz1,ny,nxp2),s(nz1)
      do 2 i =1,nx
          do 2 j =1,ny
              do 2 k =1,nz1
                  s(k) = a(k,j,i)*b(k,j,i) + s(k)
2      return
      end
      subroutine intg(a,s)
parameter(nz=32,ny=64,nx=16)
parameter(nz1=nz+1,nzml=nz-1,nzd21=nz/2+1,nzd2=nz/2)

```

```

parameter(nyd2=ny/2,nyd21=ny/2+1)
parameter(nxp2=nx+2,nxd2=nx/2,nxd21=nx/2+1)
dimension a(nyd2,nxp2,nzd21,nz1),s(nzd21)
  do 1 iz = 1,nzd21
    izp = iz
  do 1 i=1,nxd21
  do 1 j=1,nyd2
    ir = 2*i-1
1      s(iz) = a(j,ir,iz,izp) + s(iz)
  return
end
subroutine expo
parameter(ny=64,nx=16)
parameter(nxp2=nx+2,nxd21=nx/2+1)
parameter(nyt2=2*ny,nyd21=ny/2+1,nyd22=ny/2+2)
parameter(nx3p4=3*nx+4,ny5=5*ny,nxny=(nx+2)*ny)
common/work/vecx(nxp2),vecy(nyt2),vecy1(nyt2)
common/exp/exrc(nx3p4),excr(nx3p4),eycc(ny5)
  call rfft2(1,1,nx,vecx,exrc,vecx)
  call crfft2(1,1,nx,vecx,excr,vecx)
  call cfft2(1,1,ny,vecy,eycc,vecy)
return
end

```

```

program klana
cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
C   As described in the text, this code computes the
C   spatial correlation function by Fourier transformation of the
C   spectral correlations computed from the raw velocity fields
C   by using the klexp code.
C
C
cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
parameter(nz=32,ny=64,nx=16)
parameter(nym=23,nxm=14)
parameter(nz1=nz+1,nzml=nz-1,nzd21=nz/2+1,nzd2=nz/2)
parameter(nyd2=ny/2,nyd21=ny/2+1,nymt2=nym*2)
parameter(nxp2=nx+2,nxd2=mx/2,nxd21=mx/2+1,nxmd2=nxm/2)
dimension z(nz1),cnorm(nzd21),cn(3,nz1)
dimension ulu1(nym,nxm,nzd21,nz1),u2u2(nym,nxm,nzd21,nz1)
dimension u3u3(nym,nxm,nzd21,nz1),ulu2(nym,nxm,nzd21,nz1)
dimension ulu3(nym,nxm,nzd21,nz1),u2u3(nym,nxm,nzd21,nz1)
dimension phi(nym,nxm),phisym(ny,nxp2),r(nyd2,nx)
dimension rrij(nyd2,nx,nzd21,nz1)
cccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccccc
c   First read in the six components of the spectral correlation
c   function.
      do 10 izp = 1,nz1
      read(10) ((ulu1(j,i,iz,izp),j=1,nym),i=1,nxm),iz=1,nzd21)
      read(11) ((u2u2(j,i,iz,izp),j=1,nym),i=1,nxm),iz=1,nzd21)
      read(12) ((u3u3(j,i,iz,izp),j=1,nym),i=1,nxm),iz=1,nzd21)
      read(13) ((ulu2(j,i,iz,izp),j=1,nym),i=1,nxm),iz=1,nzd21)
      read(14) ((ulu3(j,i,iz,izp),j=1,nym),i=1,nxm),iz=1,nzd21)
      read(15) ((u2u3(j,i,iz,izp),j=1,nym),i=1,nxm),iz=1,nzd21)
10      continue

      DO 8000 IJ=1,6 ! loop on 11,22,33,12,13,23 & calculate rrij

      DO 5010 IZP=1,NZ1
      do 5010 iz=1,nzd21

         GOTO (5021,5022,5023,5024,5025,5026), IJ

5021      do 5031 j=1,nym
         do 5031 i=1,nxm
            phi(j,i)=ulu1(j,i,iz,izp)
5031      CONTINUE
      goto 5029

5022      do 5032 j=1,nym
         do 5032 i=1,nxm
            phi(j,i)=u2u2(j,i,iz,izp)
5032      CONTINUE
      goto 5029

5023      do 5033 j=1,nym
         do 5033 i=1,nxm
            phi(j,i)=u3u3(j,i,iz,izp)
5033      CONTINUE
      goto 5029

5024      do 5034 j=1,nym
         do 5034 i=1,nxm
            phi(j,i)=ulu2(j,i,iz,izp)

```

```

5034    CONTINUE
      goto 5029

5025    do 5035 j=1,nym
      do 5035 i=1,nxm
        phi(j,i)=ulu3(j,i,iz,izp)
5035    CONTINUE
      goto 5029

5026    do 5036 j=1,nym
      do 5036 i=1,nxm
        phi(j,i)=u2u3(j,i,iz,izp)
5036    CONTINUE

5029    continue

      call setphi(phi,phisym,IJ)
      call rij(phisym,r,cn,ij,iz,izp)

C
C NOW PUT r BACK INTO rrij - ("rij" is name of subroutine)
C
      do 5030 i=1,nx
      do 5030 j=1,nyd2
        rrij(j,i,iz,izp)=r(j,i)
5030  continue
5010  continue
6000  CONTINUE

C
C Normalize each rrij, then write it to disk:
C
      GOTO (7501,7501,7501,7504,7505,7506), IJ

7501 DO 7511 IZP=1,NZ1
      DO 7511 IZ=1,NZD21
      DO 7511 I=1,NX
      DO 7511 J=1,NYD2
        RRIJ(J,I,IZ,IZP)=RRIJ(J,I,IZ,IZP)/SQRT(CN(IJ,IZ)*CN(IJ,IZP))
7511  CONTINUE
      GOTO 7600

7504 DO 7514 IZP=1,NZ1
      DO 7514 IZ=1,NZD21
      DO 7514 I=1,NX
      DO 7514 J=1,NYD2
        RRIJ(J,I,IZ,IZP)=RRIJ(J,I,IZ,IZP)/SQRT(CN(1,IZ)*CN(2,IZP))
7514  CONTINUE
      GOTO 7600

7505 DO 7515 IZP=1,NZ1
      DO 7515 IZ=1,NZD21
      DO 7515 I=1,NX
      DO 7515 J=1,NYD2
        RRIJ(J,I,IZ,IZP)=RRIJ(J,I,IZ,IZP)/SQRT(CN(1,IZ)*CN(3,IZP))
7515  CONTINUE
      GOTO 7600

7506 DO 7516 IZP=1,NZ1
      DO 7516 IZ=1,NZD21
      DO 7516 I=1,NX
      DO 7516 J=1,NYD2

```

```

      RRIJ(J,I,IZ,IZP)=RRIJ(J,I,IZ,IZP)/SQRT(CN(2,IZ)*CN(3,IZP))
7516 CONTINUE

7600 CONTINUE
C
C DUMP rrij TO DISKFILE CONSISTENT WITH uij IN READ ROUTINE
C
    lu=80+ij-1
    write(lu) (((rrij(j,i,iz,izp),
1                           j=1,nyd2),i=1,nx),iz=1,nzd21),izp=1,nz1)

8000 CONTINUE

    write(6,110)' NORMALIZATION FACTORS:'
110 format(/a30)
    write(6,111) (((ij,izp,cn(ij,izp)),izp=1,nz1),IJ=1,3)
111 format(5(2i6,f10.6))

    stop
end

subroutine setphi(a,b,IJ)

*****
* EXPANDS uiuj OR PHIij ARRAYS FROM a(nym,nxm) TO b(ny,nxp2)
*           i. e., FROM a(23,14)   TO b(64,18)
*****
parameter (ny=64,nx=16)
parameter (nym=23,nxm=14)
parameter (nxp2=nx+2)
dimension a(nym,nxm),b(ny,nxp2)

do 1 i=1,nxm          ! 1:14
do 1 j=1,nym          ! 1:23
1   b(j,i)=a(j,i)

do 2 i=nxm+1,nxp2    ! 15:18
do 2 j=nym+1,33       ! 24:33
2   b(j,i)=0.0

do 3 i=nxm+1,nxp2    ! 15:18
do 3 j=1,nym          ! 1:23
3   b(j,i)=0.0

do 4 i=1,nxm          ! 1,14
do 4 j=nym+1,33       ! 24:33
4   b(j,i)=0.0

IF (IJ.LE.4) THEN      ! 11,22,33,12
do 10 i=1,nxp2         ! 1:18
do 10 j=34,ny           ! 34:64
10   b(j,i)= b(ny-j+2,i)
ENDIF

IF (IJ.GE.5) THEN      ! 13,23 have odd symmetry!
do 11 i=1,nxp2         ! 1:18
do 11 j=34,ny           ! 34:64
11   b(j,i)=-b(ny-j+2,i)
ENDIF

```

```

return
end

subroutine rij(a,b,cn,ij,iz,izp)

*****  

* INPUTS EXPANDED ARRAY FROM SETPHI, OUTPUTS CORRELATION FUNCTN. Rij *  

*****  

parameter (nz=32,nz1=nz+1,ny=64,nx=16)
parameter (nxp2=nx+2,nxd21=nx/2+1)
parameter (nyt2=2*ny,nyd21=ny/2+1)
parameter (nyd2=ny/2,nxd2=nx/2,ny5=5*ny,nx3p4=3*nx+4)
dimension eyy(ny5),exx(nx3p4)
dimension vecy(nyt2),vecx(nxp2)
dimension a(ny,nxp2),b(nyd2,nxp2),c(nyd2,nxp2),cn(3,nz1)

c
c INITIALIZE FOURIER TRANSFORMS (CALC. eyy AND exx):
c
call cfft2 (1,1,ny,vecy,eyy,vecy)
call crfft2(1,1,nx,vecx,exx,vecx)
c
c EXTRACT VECTOR FROM ARRAY AND DO FFT IN Y DIRECTION:
c
do 1 i=1,nxd21

do 2 j=1,ny
  vecy(2*j-1)=a(j,2*i-1)      ! REAL PART
2  vecy(2*j )=a(j,2*i )      ! IMAG PART
  call cfft2(0,1,ny,vecy,eyy,vecy)

do 3 j=1,nyd2                  ! PUT VECY INTO TEMPORARY OUTPUT
  c(j,2*i-1)=vecy(2*j-1)
3  c(j,2*i )=vecy(2*j )

1 continue

c
c EXTRACT VECTOR FROM ARRAY AND DO FFT IN X DIRECTION:
c
do 4 j=1,nyd2

do 5 i=1,nxp2
5  vecx(i)=c(j,i)
  call crfft2(0,1,nx,vecx,exx,vecx)

do 6 i=1,nx                    ! PUT VECX INTO FINAL OUTPUT
6  c(j,i)=vecx(i)

4 continue

if (IJ.LE.3.AND.izp.eq.iz) then      ! set normalization factors
  cn(ij,iz)=c(1,1)                  ! ij=1,3 for 11,22,33
  if (cn(ij,iz).eq.0.) cn(ij,iz)=1.
  cn(ij,nz1+1-iz)=cn(ij,iz)
endif

c

```

```
c FINALLY, FLIP RIGHT SIDE TO THE LEFT OF THE LEFT SIDE:  
c  
    do 10 j=1,nyd2  
    do 11 i=1,nxd2  
        b(j,i )=c(j,i+8)  
11      b(j,i+8)=c(j,i )  
10  continue  
  
    return  
end
```

```

program wavspec

*****
*      "Computes space-time corelations for planar data"          *
* THIS VERSION READS 7 FILES, DOES 5 FFT'S, AND AVOIDS ZERO PACKING   *
*****
* This version reads three 3-d yxt files u1, u2, and u3, extracts from*
* each three 2-d arrays (umn,unk,umk), does 2-d fft on each of those   *
* still using same array names, and adds up the results over the     *
* remaining (third) variable. It then calculates energies, ensemble   *
* averages, and computes correlation fundtions, saving all           *
* results to disk.                                                 *
*****
parameter (ny=64,nx=16,ntim=512)
dimension u1(ny,nx,ntim),u2(ny,nx,ntim),u3(ny,nx,ntim)
dimension umn1(ny , nx+2),umn2(ny , nx+2),umn3(ny , nx+2)
dimension smn1(ny/2,nx/2+1),smn2(ny/2,nx/2+1),smn3(ny/2,nx/2+1)
dimension smn1s(ny , nx+2),smn2s(ny , nx+2),smn3s(ny , nx+2)
dimension rmn1(ny/2, nx+2),rmn2(ny/2, nx+2),rmn3(ny/2, nx+2)
dimension unk1(nx , ntim+2),unk2(nx , ntim+2),unk3(nx, ntim+2)
dimension snk1(nx,ntim/2+1),snk2(nx,ntim/2+1),snk3(nx,ntim/2+1)
dimension snk1s(nx, ntim+2),snk2s(nx, ntim+2),snk3s(nx, ntim+2)
dimension rnk1(nx/2,ntim+2),rnk2(nx/2,ntim+2),rnk3(nx/2,ntim+2)
dimension umk1(ny,ntim+2) ,umk2(ny, ntim+2),umk3(ny, ntim+2)
dimension smk1(ny/2,ntim/2+1),smk2(ny/2,ntim/2+1),
*          smk3(ny/2,ntim/2+1)
dimension smk1s(ny,ntim+2) ,smk2s(ny,ntim+2) ,smk3s(ny,ntim+2)
dimension rmk1(ny/2,ntim+2),rmk2(ny/2,ntim+2),rmk3(ny/2,ntim+2)
dimension tmp(ny,nx),e(9),eavg(3),w(9),tmp1(nx,ntim/2+1)

c  initialize 9 ensemble averages smni, ski, smki, i=1,3 :
c
do 40 m=1,ny/2
do 40 n=1,nx/2+1
  smn1(m,n)=0.
  smn2(m,n)=0.
  smn3(m,n)=0.
40 continue
do 43 n=1,nx
do 43 k=1,ntim/2+1
  snk1(n,k)=0.
  snk2(n,k)=0.
  snk3(n,k)=0.
43 continue
do 46 m=1,ny/2
do 46 k=1,ntim/2+1
  smk1(m,k)=0.
  smk2(m,k)=0.
  smk3(m,k)=0.
46 continue

c  7 files (lu=10,16) permit 5 FFT's with 240 k's left over at the end
c
call read  (u1,u2,u3,tmp,nx,ny,ntim,10, 1,400) ! 400
call read  (u1,u2,u3,tmp,nx,ny,ntim,11,401,512) ! 112
call hfens (u1,u2,u3,umn1,umn2,umn3,smn1,smn2,smn3,unk1,unk2,unk3,
* snk1,snk2,snk3,umk1,umk2,umk3,smk1,smk2,smk3,nx,ny,ntim) ! ====
call read  (u1,u2,u3,tmp,nx,ny,ntim,11, 1,288) ! 288
call read  (u1,u2,u3,tmp,nx,ny,ntim,12,289,512) ! 224

```

```

call hfens (u1,u2,u3,umn1,umn2,umn3,smn1,smn2,smn3,unk1,unk2,unk3,
* snk1,snk2,snk3,umk1,umk2,umk3,smk1,smk2,smk3,nx,ny,ntim) !=====
call read (u1,u2,u3,tmp(nx,ny,ntim,12, 1,176) ! 176
call read (u1,u2,u3,tmp(nx,ny,ntim,13,177,512) ! 336
call hfens (u1,u2,u3,umn1,umn2,umn3,smn1,smn2,smn3,unk1,unk2,unk3,
* snk1,snk2,snk3,umk1,umk2,umk3,smk1,smk2,smk3,nx,ny,ntim) !=====
call read (u1,u2,u3,tmp(nx,ny,ntim,13, 1, 64) ! 64
call read (u1,u2,u3,tmp(nx,ny,ntim,14, 65,464) ! 400
call read (u1,u2,u3,tmp(nx,ny,ntim,15,465,512) ! 48
call hfens (u1,u2,u3,umn1,umn2,umn3,smn1,smn2,smn3,unk1,unk2,unk3,
* snk1,snk2,snk3,umk1,umk2,umk3,smk1,smk2,smk3,nx,ny,ntim) !=====
call read (u1,u2,u3,tmp(nx,ny,ntim,15, 1,352) ! 352
call read (u1,u2,u3,tmp(nx,ny,ntim,16,353,512) ! 160
call hfens (u1,u2,u3,umn1,umn2,umn3,smn1,smn2,smn3,unk1,unk2,unk3,
* snk1,snk2,snk3,umk1,umk2,umk3,smk1,smk2,smk3,nx,ny,ntim) !=====

c
c calculate and print energies:
c

call energ1(smn1, ny, nx, e(1),w(1))
call energ1(smn2, ny, nx, e(2),w(2))
call energ1(smn3, ny, nx, e(3),w(3))
e(1)=sqrt(e(1)/5./512.) !
e(2)=sqrt(e(2)/5./512.) ! average over 5*512 data pts
e(3)=sqrt(e(3)/5./512.) !
print *, ' MN RMS values 1, 2, 3 are: ',(e(i),i=1,3)

call energ2(snk1, nx, ntim, e(4),w(4))
call energ2(snk2, nx, ntim, e(5),w(5))
call energ2(snk3, nx, ntim, e(6),w(6))
e(4)=sqrt(e(4)/64.*512./512.) !
e(5)=sqrt(e(5)/64.*512./512.) ! energ2 divides by 512**2
e(6)=sqrt(e(6)/64.*512./512.) !
print *, ' NK RMS values 4, 5., 6 are: ',(e(i),i=4,6)

call energ3(smk1, ny, ntim, e(7),w(7))
call energ3(smk2, ny, ntim, e(8),w(8))
call energ3(smk3, ny, ntim, e(9),w(9))
e(7)=sqrt(e(7)/16.*512./512.) !
e(8)=sqrt(e(8)/16.*512./512.) ! energ3 divides by 512**2
e(9)=sqrt(e(9)/16.*512./512.) !
print *, ' MK RMS values 7, 8, 9 are: ',(e(i),i=7,9)

c
c test energies:
c

do 1050 i=1,3
  eavg(i)=0.
  do 1010 j=i,9,3
    eavg(i)=eavg(i)+e(j)
1010  continue
  eavg(i)=eavg(i)/3.
  do 1020 j=i,9,3
    if(abs(e(j)-eavg(i))/eavg(i).gt.0.00001) then
      print *, ' energy error! e(' ,j,') = ',e(j),'; avg = ',eavg(i)
      endif
1020  continue
1050 continue

c
c normalize ensemble averages by energy:
c

do 1100 m=1,ny/2

```

```

      do 1100 n=1,nx/2+1
        smn1(m,n)=smn1(m,n)/w(1)
        smn2(m,n)=smn2(m,n)/w(2)
        smn3(m,n)=smn3(m,n)/w(3)
1100 continue
      do 1150 n=1,nx
        do 1150 k=1,ntim/2+1
          snk1(n,k)=snk1(n,k)/w(4)
          snk2(n,k)=snk2(n,k)/w(5)
          snk3(n,k)=snk3(n,k)/w(6)
1150 continue
      do 1200 m=1,ny/2
        do 1200 k=1,ntim/2+1
          smk1(m,k)=smk1(m,k)/w(7)
          smk2(m,k)=smk2(m,k)/w(8)
          smk3(m,k)=smk3(m,k)/w(9)
1200 continue
c
c  symmetrize smni, snki, smki:
c
      call setphi1(smn1,smn1s,ny, nx)
      call setphi1(smn2,smn2s,ny, nx)
      call setphi1(smn3,smn3s,ny, nx)
      call setphi2(snk1,snk1s,nx,ntim)
      call setphi2(snk2,snk2s,nx,ntim)
      call setphi2(snk3,snk3s,nx,ntim)
      call setphi1(smk1,smk1s,ny,ntim)
      call setphi1(smk2,smk2s,ny,ntim)
      call setphi1(smk3,smk3s,ny,ntim)
c
c  calculate correlation functions:
c
      call rij1(smn1s,rmn1,ny, nx)
      call rij1(smn2s,rmn2,ny, nx)
      call rij1(smn3s,rmn3,ny, nx)
      call rij2(snk1s,rnk1,nx,ntim)
      call rij2(snk2s,rnk2,nx,ntim)
      call rij2(snk3s,rnk3,nx,ntim)
      call rij3(smk1s,rmk1,ny,ntim)
      call rij3(smk2s,rmk2,ny,ntim)
      call rij3(smk3s,rmk3,ny,ntim)
c
c  flip snk spectra now that fft's are completed:
c
      do 1310 n=1,nx
        do 1310 k=1,ntim/2+1
          tmp1(n,k)=snk1(n,k)
1310 continue
      do 1320 n=1,nx/2
        do 1320 k=1,ntim/2+1
          snk1(n,k)=tmp1(n+nx/2,k)
1320 continue
      do 1330 n=nx/2+1,nx
        do 1330 k=1,ntim/2+1
          snk1(n,k)=tmp1(n-nx/2,k)
1330 continue
      do 1410 n=1,nx
        do 1410 k=1,ntim/2+1
          tmp1(n,k)=snk2(n,k)
1410 continue

```

```

        do 1420 n=1,nx/2
        do 1420 k=1,ntim/2+1
          snk2(n,k)=tmp1(n+nx/2,k)
1420 continue
        do 1430 n=nx/2+1,nx
        do 1430 k=1,ntim/2+1
          snk2(n,k)=tmp1(n-nx/2,k)
1430 continue
        do 1510 n=1,nx
        do 1510 k=1,ntim/2+1
          tmp1(n,k)=snk3(n,k)
1510 continue
        do 1520 n=1,nx/2
        do 1520 k=1,ntim/2+1
          snk3(n,k)=tmp1(n+nx/2,k)
1520 continue
        do 1530 n=nx/2+1,nx
        do 1530 k=1,ntim/2+1
          snk3(n,k)=tmp1(n-nx/2,k)
1530 continue
c
c  save flipped (where applicable) spectra to disk:
c
      write(50,91) 1,e(1),eavg(1)
      write(50,92) ((smn1(m,n), m=1,ny/2), n=1, nx/2+1)
      write(50,91) 2,e(2),eavg(2)
      write(50,92) ((smn2(m,n), m=1,ny/2), n=1, nx/2+1)
      write(50,91) 3,e(3),eavg(3)
      write(50,92) ((smn3(m,n), m=1,ny/2), n=1, nx/2+1)
      write(50,91) 4,e(4),eavg(1)
      write(50,92) ((snk1(n,k), n=1, nx), k=1,ntim/2+1)
      write(50,91) 5,e(5),eavg(2)
      write(50,92) ((snk2(n,k), n=1, nx), k=1,ntim/2+1)
      write(50,91) 6,e(6),eavg(3)
      write(50,92) ((snk3(n,k), n=1, nx), k=1,ntim/2+1)
      write(50,91) 7,e(7),eavg(1)
      write(50,92) ((smk1(m,k), m=1,ny/2), k=1,ntim/2+1)
      write(50,91) 8,e(8),eavg(2)
      write(50,92) ((smk2(m,k), m=1,ny/2), k=1,ntim/2+1)
      write(50,91) 9,e(9),eavg(3)
      write(50,92) ((smk3(m,k), m=1,ny/2), k=1,ntim/2+1)

      91 format(2x,i3,2e12.5)
      92 format(2x,4e12.5)
c
c  save correlation functions:
c
      write(50,91) 1
      write(50,92) ((rmn1(m,n),m=1,ny/2),n=1, nx)
      write(50,91) 2
      write(50,92) ((rmn2(m,n),m=1,ny/2),n=1, nx)
      write(50,91) 3
      write(50,92) ((rmn3(m,n),m=1,ny/2),n=1, nx)
      write(50,91) 4
      write(50,92) ((rnk1(n,k),n=1,nx/2),k=1,ntim)
      write(50,91) 5
      write(50,92) ((rnk2(n,k),n=1,nx/2),k=1,ntim)
      write(50,91) 6
      write(50,92) ((rnk3(n,k),n=1,nx/2),k=1,ntim)
      write(50,91) 7

```

```

write(50,92) ((rmk1(m,k),m=1,ny/2),k=1,ntim)
write(50,91) 8
write(50,92) ((rmk2(m,k),m=1,ny/2),k=1,ntim)
write(50,91) 9
write(50,92) ((rmk3(m,k),m=1,ny/2),k=1,ntim)

stop
end

subroutine read(u1,u2,u3,tmp,nx,ny,ntim,lu,ka,kz)
dimension u1(ny,nx,ntim),u2(ny,nx,ntim),u3(ny,nx,ntim),tmp(ny,nx)
do 50 k=ka,kz
    read (lu) ((u1(j,i,k),j=1,ny),i=1,nx)
    read (lu) ((u3(j,i,k),j=1,ny),i=1,nx)
    read (lu) ((u2(j,i,k),j=1,ny),i=1,nx)
    read (lu) ((tmp(j,i),j=1,ny),i=1,nx)
    read (lu) ((tmp(j,i),j=1,ny),i=1,nx)
    read (lu) ((tmp(j,i),j=1,ny),i=1,nx)
50 continue
return
end

subroutine hfens (u1,u2,u3,      umn1,umn2,umn3,smn1,smn2,smn3,
*   unk1,unk2,unk3,snk1,snk2,snk3,umk1,umk2,umk3,smk1,smk2,smk3,
*   nx,ny,ntim)

dimension u1(ny,nx,ntim),u2(ny,nx,ntim),u3(ny,nx,ntim)
dimension umn1(ny , nx+2),umn2(ny , nx+2),umn3(ny , nx+2)
dimension smn1(ny/2,nx/2+1),smn2(ny/2,nx/2+1),smn3(ny/2,nx/2+1)
dimension unk1(nx , ntim+2),unk2(nx , ntim+2),unk3(nx , ntim+2)
dimension snk1(nx,ntim/2+1),snk2(nx,ntim/2+1),snk3(nx,ntim/2+1)
dimension umk1(ny,ntim+2) ,umk2(ny, ntim+2),umk3(ny, ntim+2)
dimension smk1(ny/2,ntim/2+1),smk2(ny/2,ntim/2+1),
*           smk3(ny/2,ntim/2+1)

c   At each time (3rd variable), extract umn array, fft it, compile sum:
c
do 100 k=1,ntim
    do 70 m=1,ny
        do 70 n=1,nx
            umn1(m,n)=u1(m,n,k)
            umn2(m,n)=u2(m,n,k)
            umn3(m,n)=u3(m,n,k)
70     continue
        call hfft1(umn1,ny,nx,1)
        call ensav1(umn1,smn1,ny,nx)
        call hfft1(umn2,ny,nx,1)
        call ensav1(umn2,smn2,ny,nx)
        call hfft1(umn3,ny,nx,1)
        call ensav1(umn3,smn3,ny,nx)
100    continue
c   At each y (3rd variable), extract unk array, fft it, compile sum:
c
    do 200 m=1,ny

```

```

do 170 n=1,nx
do 170 k=1,ntim
  unk1(n,k)=u1(m,n,k)
  unk2(n,k)=u2(m,n,k)
  unk3(n,k)=u3(m,n,k)
170 continue
call hfft2(unk1,nx,ntim,1)
call ensav2(unk1,snk1,nx,ntim)
call hfft2(unk2,nx,ntim,1)
call ensav2(unk2,snk2,nx,ntim)
call hfft2(unk3,nx,ntim,1)
call ensav2(unk3,snk3,nx,ntim)
200 continue
c
c At each x (3rd variable), extract umk array, fft it, compile sum:
c
do 300 n=1,nx
  do 270 m=1,ny
    do 270 k=1,ntim
      umk1(m,k)=u1(m,n,k)
      umk2(m,k)=u2(m,n,k)
      umk3(m,k)=u3(m,n,k)
270 continue
call hfft3(umk1,ny,ntim,1)
call ensav3(umk1,smk1,ny,ntim)
call hfft3(umk2,ny,ntim,1)
call ensav3(umk2,smk2,ny,ntim)
call hfft3(umk3,ny,ntim,1)
call ensav3(umk3,smk3,ny,ntim)
300 continue
return
end

```

```

subroutine hfft1(a,n1,n2,is)
parameter (m1=64,m2=16) ! needed to dimension noncalled arrays
dimension a(n1,n2+2)
dimension indx(m2+2),indy(2*m1)
dimension exx(3*m2+4),exy(5*m1)
dimension vecx(m2+2),vecy(2*m1),vecyl(2*m1)
  if (is.eq. -1) go to 500
call rcfft2(1,1,n2,vecx,exx,vecx)
call cfft2(1,1,n1,vecy,exy,vecy)
do 1 j=1,n1
  do 2 i=1,n2+2
    indx(i)=(i-1)*n1+j
2 continue
  call gather(n2,vecx,a,indx)
  call rcfft2(0,1,n2,vecx,exx,vecx)
  call scatter(n2+2,a,indx,vecx)
1 continue
  do 3 i=1,n2/2+1
  do 4 j=1,n1
    j1=2*(j-1)+1
    indy(j1)=j+2*(i-1)*n1
    indy(j1+1)=indy(j1)+n1
4 continue
  call gather(2*n1,vecy,a,indy)
  call cfft2(0,1,n1,vecy,exy,vecy)

```

```

    call scatter(2*n1,a,indy,vecy)
3  continue
    return
500 continue
    call crfft2(1,-1,n2,vecx,exx,vecx)
    call cfft2(1,-1,n1,vecy,exy,vecy)
    do 5 i=1,n2/2+1
    do 6 j=1,n1
        j1=2*(j-1)+1
        indy(j1)=j+2*(i-1)*n1
        indy(j1+1)=indy(j1)+n1
6  continue
    call gather(2*n1,vecy,a,indy)
    call cfft2(0,-1,n1,vecy,exy,vecy)
    call scatter(2*n1,a,indy,vecy)
5  continue
    do 7 j=1,n1
    do 8 i=1,n2+2
        indx(i)=(i-1)*n1+j
8  continue
    call gather(n2+2,vecx,a,indx)
    call crfft2(0,-1,n2,vecx,exx,vecx)
    call scatter(n2,a,indx,vecx)
7  continue
    do 10 i=1,n2
    do 10 j=1,n1
        a(j,i)=a(j,i)/(2.*float(n2*n1))
10 continue
    return
end

subroutine hfft2(a,n1,n2,is)
parameter (m1=16,m2=512) ! needed to dimension noncalled arrays
dimension a(n1,n2+2)
dimension indx(m2+2),indy(2*m1)
dimension exx(3*m2+4),exy(5*m1)
dimension vecx(m2+2),vecy(2*m1),vecy1(2*m1)
    if (is.eq. -1) go to 500
call rcfft2(1,1,n2,vecx,exx,vecx)
call cfft2(1,1,n1,vecy,exy,vecy)
do 1 j=1,n1
do 2 i=1,n2+2
    indx(i)=(i-1)*n1+j
2  continue
    call gather(n2,vecx,a,indx)
    call rcfft2(0,1,n2,vecx,exx,vecx)
    call scatter(n2+2,a,indx,vecx)
1  continue
    do 3 i=1,n2/2+1
    do 4 j=1,n1
        j1=2*(j-1)+1
        indy(j1)=j+2*(i-1)*n1
        indy(j1+1)=indy(j1)+n1
4  continue
    call gather(2*n1,vecy,a,indy)
    call cfft2(0,1,n1,vecy,exy,vecy)
    call scatter(2*n1,a,indy,vecy)
3  continue
    return

```

```

500 continue
call crfft2(1,-1,n2,vecx,exx,vecx)
call cfft2(1,-1,n1,vecy,exy,vecy)
do 5 i=1,n2/2+1
do 6 j=1,n1
    j1=2*(j-1)+1
    indy(j1)=j+2*(i-1)*n1
    indy(j1+1)=indy(j1)+n1
6   continue
    call gather(2*n1,vecy,a,indy)
    call cfft2(0,-1,n1,vecy,exy,vecy)
    call scatter(2*n1,a,indy,vecy)
5   continue
    do 7 j=1,n1
    do 8 i=1,n2+2
        indx(i)=(i-1)*n1+j
8   continue
    call gather(n2+2,vecx,a,indx)
    call crfft2(0,-1,n2,vecx,exx,vecx)
    call scatter(n2,a,indx,vecx)
7   continue
    do 10 i=1,n2
    do 10 j=1,n1
        a(j,i)=a(j,i)/(2.*float(n2*n1))
10  continue
return
end

subroutine hfft3(a,n1,n2,is)
parameter (m1=64,m2=512) ! needed to dimension noncalled arrays
dimension a(n1,n2+2)
dimension indx(m2+2),indy(2*m1)
dimension exx(3*m2+4),exy(5*m1)
dimension vecx(m2+2),vecy(2*m1),vecy1(2*m1)
    if (is.eq. -1) go to 500
call rcffft2(1,1,n2,vecx,exx,vecx)
call cfft2(1,1,n1,vecy,exy,vecy)
do 1 j=1,n1
do 2 i=1,n2+2
    indx(i)=(i-1)*n1+j
2   continue
    call gather(n2,vecx,a,indx)
    call rcffft2(0,1,n2,vecx,exx,vecx)
    call scatter(n2+2,a,indx,vecx)
1   continue
    do 3 i=1,n2/2+1
    do 4 j=1,n1
        j1=2*(j-1)+1
        indy(j1)=j+2*(i-1)*n1
        indy(j1+1)=indy(j1)+n1
4   continue
    call gather(2*n1,vecy,a,indy)
    call cfft2(0,1,n1,vecy,exy,vecy)
    call scatter(2*n1,a,indy,vecy)
3   continue
return
500 continue
call crfft2(1,-1,n2,vecx,exx,vecx)
call cfft2(1,-1,n1,vecy,exy,vecy)

```

```

do 5 i=1,n2/2+1
do 6 j=1,n1
    j1=2*(j-1)+1
    indy(j1)=j+2*(i-1)*n1
    indy(j1+1)=indy(j1)+n1
6 continue
call gather(2*n1,vecy,a,indy)
call cfft2(0,-1,n1,vecy,exy,vecy)
call scatter(2*n1,a,indy,vecy)
5 continue
do 7 j=1,n1
do 8 i=1,n2+2
    indx(i)=(i-1)*n1+j
8 continue
    call gather(n2+2,vecx,a,indx)
    call crfft2(0,-1,n2,vecx,exx,vecx)
    call scatter(n2,a,indx,vecx)
7 continue
do 10 i=1,n2
do 10 j=1,n1
    a(j,i)=a(j,i)/(2.*float(n2*n1))
10 continue
return
end

subroutine ensav1(a,b,n1,n2)
parameter (m1=64,m2=16)
dimension a(n1,n2+2),b(n1/2,n2/2+1)
dimension r(m1,m2/2+1) ! need var dim = r(n1,n2/2+1)
do 1 i=1,n2/2+1
do 1 j=1,n1
    r(j,i)=a(j,2*i-1)**2+a(j,2*i)**2
1 continue
call sym1(r,n1,n2/2+1)
do 2 i=1,n2/2+1
do 2 j=1,n1/2
    b(j,i)=b(j,i)+r(j,i)
2 continue
return
end

subroutine ensav2(a,b,n1,n2)
dimension a(n1,n2+2),b(n1,n2/2+1)
do 2 i=1,n2/2+1
do 2 j=1,n1
    b(j,i)=b(j,i)+a(j,2*i-1)**2+a(j,2*i)**2
2 continue
return
end

subroutine ensav3(a,b,n1,n2)
parameter (m1=64,m2=512)
dimension a(n1,n2+2),b(n1/2,n2/2+1)
dimension r(m1,m2/2+1) ! need var dim = r(n1,n2/2+1)
do 1 i=1,n2/2+1
do 1 j=1,n1
    r(j,i)=a(j,2*i-1)**2+a(j,2*i)**2

```

```

1      continue
2      call sym1(r,n1,n2/2+1)
3      do 2 i=1,n2/2+1
4      do 2 j=1,n1/2
5          b(j,i)=b(j,i)+r(j,i)
6      continue
7      return
8      end

subroutine sym1(c,m1,m2)
dimension c(m1,m2)
do 1 i=1,m2
do 1 j=1,m1/2
    if ( j .eq. 1) then
        c(j,i) = c(j,i)
    else
        c(j,i) = (c(j,i) + c(m1-j+2,i))/2.
    endif
1  continue
return
end

subroutine energ1(b,n1,n2,e,w)
parameter (m1=64,m2=16) ! needed to dimension noncalled arrays
dimension b(n1/2,n2/2+1)
dimension c(m1/2,m2/2+1)
do 4 j=2,n1/2
    c(j,1)=b(j,1)/2.
    b(j,1)=b(j,1)/8.
4  continue
do 5 i=2,n2/2+1
    c(1,i)=b(1,i)/2.
    b(1,i)=b(1,i)/8.
5  continue
    c(1,1)=0.
    b(1,1)=b(1,1)/16.
    do 6 i=2,n2/2+1
        do 6 j=2,n1/2
            c(j,i)=b(j,i)
            b(j,i)=b(j,i)/4.
6  continue
e=0.
w=0.
do 7 i=1,n2/2+1
do 7 j=1,n1/2
    w=w+c(j,i)
    e=e+b(j,i)
7  continue
e=e*4./n1/n1/n2/n2
return
end

subroutine energ2(b,n1,n2,e,w)
parameter (m1=16,m2=512) ! needed to dimension noncalled arrays
dimension b(n1,n2/2+1)
dimension c(m1,m2/2+1)
do 4 j=2,n1

```

```

        c(j,1)=b(j,1)/4.
        b(j,1)=b(j,1)/16.
4 continue
    do 5 i=2,n2/2+1
        c(1,i)=b(1,i)/2.
        b(1,i)=b(1,i)/8.
5 continue
    c(1,1)=0.
    b(1,1)=b(1,1)/16.
    do 6 i=2,n2/2+1
    do 6 j=2,n1
        c(j,i)=b(j,i)/2.
        b(j,i)=b(j,i)/8.
6 continue
w=0.
e=0.
    do 7 i=1,n2/2+1
    do 7 j=1,n1
        w=w+c(j,i)
        e=e+b(j,i)
7 continue
e=e*4./n1/n1/n2/n2
return
end

subroutine energ3(b,n1,n2,e,w)
parameter (m1=64,m2=512) ! needed to dimension noncalled arrays
dimension b(n1/2,n2/2+1)
dimension c(m1/2,m2/2+1)
do 4 j=2,n1/2
    c(j,1)=b(j,1)/2.
    b(j,1)=b(j,1)/8.
4 continue
    do 5 i=2,n2/2+1
        c(1,i)=b(1,i)/2.
        b(1,i)=b(1,i)/8.
5 continue
    c(1,1)=0.
    b(1,1)=b(1,1)/16.
    do 6 i=2,n2/2+1
    do 6 j=2,n1/2
        c(j,i)=b(j,i)
        b(j,i)=b(j,i)/4.
6 continue
e=0.
w=0.
    do 7 i=1,n2/2+1
    do 7 j=1,n1/2
        w=w+c(j,i)
        e=e+b(j,i)
7 continue
e=e*4./n1/n1/n2/n2
return
end

subroutine setphi1(a,b,n1,n2)
*****

```

```

* EXPANDS smni    ARRAYS   FROM a(ny/2,nx/2+1)      TO b(ny,nx+2)      *
*                      i. e.,   FROM a(32,          9)      TO b(64, 18)      *
*                     smki    ARRAYS   FROM a(ny/2,ntim/2+1)  TO b(ny,ntim+2)  *
*                      i. e.,   FROM a(32,        257)      TO b(64,514)      *
*****
```

```

dimension a(n1/2,n2/2+1),b(n1,n2+2)

do 1 i=1,n2/2+1      ! 1: 9,257
do 1 j=1,n1/2        ! 1: 32, 32
  b(j,2*i-1) =a(j,i) ! a into real(b)
  b(j,2*i ) =0.       ! imag(b)=0.
1 continue

do 2 i=1,n2+2        ! 1: 18,514
  b(n1/2+1,i)=0.     ! row 33, 33 = 0.
2 continue

do 10 i=1,n2+2       ! 1: 10, 514
do 10 j=n1/2+2,n1    ! 34:64, 34:64
  b(j,i)= b(n1-j+2,i)! symmetry for smni, smki
10 continue
b(1,1)=0.             ! eliminate (0,0) mode

return
end
```

```
subroutine setphi2(a,b,n1,n2)
```

```

*****
```

```

* EXPANDS snki    ARRAYS   FROM a(nx, ntim/2+1)  TO b(nx, ntim+2)      *
*                      i. e.,   FROM a(16, 257)        TO b(16, 514 )      *
*****
```

```

dimension a(n1,n2/2+1),b(n1,n2+2)

do 1 i=1,n2/2+1      ! 1: 257
do 1 j=1,n1           ! 1: 16
  b(j,2*i-1) =a(j,i) ! a into real(b)
  b(j,2*i ) =0.       ! imag(b)=0.
1 continue
b(1,1)=0.             ! eliminate (0,0) mode

return
end
```

```
subroutine rij1(a,b,n1,n2)
```

```

*****
```

```

* INPUTS EXPANDED ARRAY FROM SETPHI, OUTPUTS CORRELATION FUNCTN. Rij *
*****
```

```

parameter (m1=64,m2=16) ! needed for uncalled array dimensions
dimension eyy(5*m1),exx(3*m2+4)
dimension vecy(2*m1),vecx(m2+2)
dimension a(n1,n2+2),b(n1/2,n2+2),c(m1/2,m2+2)
```

```
c
c INITIALIZE FOURIER TRANSFORMS (CALC. eyy AND exx):
```

```

c
    call cfft2 (1,1,n1,vecy,eyy,vecy)
    call crfft2(1,1,n2,vecx,exx,vecx)
c
c EXTRACT VECTOR FROM ARRAY AND DO FFT IN Y DIRECTION:
c
    do 1 i=1,n2/2+1

        do 2 j=1,n1
            vecy(2*j-1)=a(j,2*i-1)      ! REAL PART
2        vecy(2*j  )=a(j,2*i  )      ! IMAG PART
            call cfft2(0,1,n1,vecy,eyy,vecy)

        do 3 j=1,n1/2                  ! PUT VECY INTO TEMPORARY OUTPUT
            c(j,2*i-1)=vecy(2*j-1)
3        c(j,2*i  )=vecy(2*j  )

1    continue

c
c EXTRACT VECTOR FROM ARRAY AND DO FFT IN X DIRECTION:
c
    do 4 j=1,n1/2

        do 5 i=1,n2+2
5        vecx(i)=c(j,i)
            call crfft2(0,1,n2,vecx,exx,vecx)

        do 6 i=1,n2                  ! PUT VECX INTO FINAL OUTPUT
6        c(j,i)=vecx(i)

4    continue

    cnorm=c(1,1)
    if (cnorm.eq.0.) cnorm=1.

c
c FINALLY, FLIP RIGHT SIDE TO THE LEFT OF THE LEFT SIDE:
c
    do 10 j=1,n1/2
    do 11 i=1,n2/2
        b(j,i     )=c(j,i+n2/2)/cnorm
11        b(j,i+n2/2)=c(j,i     )/cnorm
10    continue

    return
end

subroutine rij2(a,b,n1,n2)

*****
* INPUTS EXPANDED ARRAY FROM SETPHI, OUTPUTS CORRELATION FUNCTN. Rij *
*****

parameter (m1=16,m2=512) ! needed for uncalled array dimensions
dimension eyy(5*m1),exx(3*m2+4)
dimension vecy(2*m1),vecx(m2+2)
dimension a(n1,n2+2),b(n1/2,n2+2),c(m1/2,m2+2)
c
c INITIALIZE FOURIER TRANSFORMS (CALC. eyy AND exx):
c

```

```

call cfft2 (1,1,n1,vecy,eyy,vecy)
call crfft2(1,1,n2,vecx,exx,vecx)
c
c EXTRACT VECTOR FROM ARRAY AND DO FFT IN Y DIRECTION:
c
do 1 i=1,n2/2+1

do 2 j=1,n1
  vecy(2*j-1)=a(j,2*i-1)      ! REAL PART
2  vecy(2*j    )=a(j,2*i     )    ! IMAG PART
  call cfft2(0,1,n1,vecy,eyy,vecy)

do 3 j=1,n1/2                  ! PUT VECY INTO TEMPORARY OUTPUT
  c(j,2*i-1)=vecy(2*j-1)
3  c(j,2*i    )=vecy(2*j    )

1 continue

c
c EXTRACT VECTOR FROM ARRAY AND DO FFT IN X DIRECTION:
c
do 4 j=1,n1/2

do 5 i=1,n2+2
5  vecx(i)=c(j,i)
  call crfft2(0,1,n2,vecx,exx,vecx)

do 6 i=1,n2                  ! PUT VECX INTO FINAL OUTPUT
6  c(j,i)=vecx(i)

4 continue

cnorm=c(1,1)
if (cnorm.eq.0.) cnorm=1.

c
c FINALLY, FLIP RIGHT SIDE TO THE LEFT OF THE LEFT SIDE:
c
do 10 j=1,n1/2
  do 11 i=1,n2/2
    b(j,i        )=c(j,i+n2/2)/cnorm
11  b(j,i+n2/2)=c(j,i        )/cnorm
10 continue

return
end

subroutine rij3(a,b,n1,n2)

*****
* INPUTS EXPANDED ARRAY FROM SETPHI, OUTPUTS CORRELATION FUNCTN. Rij *
*****

parameter (m1=64,m2=512) ! needed for uncalled array dimensions
dimension eyy(5*m1),exx(3*m2+4)
dimension vecy(2*m1),vecx(m2+2)
dimension a(n1,n2+2),b(n1/2,n2+2),c(m1/2,m2+2)
c
c INITIALIZE FOURIER TRANSFORMS (CALC. eyy AND exx):
c
call cfft2 (1,1,n1,vecy,eyy,vecy)

```

```

    call crfft2(1,1,n2,vecx,exx,vecx)
c
c EXTRACT VECTOR FROM ARRAY AND DO FFT IN Y DIRECTION:
c
    do 1 i=1,n2/2+1

        do 2 j=1,n1
            vecy(2*j-1)=a(j,2*i-1)      ! REAL PART
2        vecy(2*j   )=a(j,2*i   )      ! IMAG PART
            call cfft2(0,1,n1,vecy,eyy,vecy)

        do 3 j=1,n1/2                  ! PUT VECY INTO TEMPORARY OUTPUT
            c(j,2*i-1)=vecy(2*j-1)
3        c(j,2*i   )=vecy(2*j   )

    1 continue

c
c EXTRACT VECTOR FROM ARRAY AND DO FFT IN X DIRECTION:
c
    do 4 j=1,n1/2

        do 5 i=1,n2+2
5        vecx(i)=c(j,i)
            call crfft2(0,1,n2,vecx,exx,vecx)

        do 6 i=1,n2                  ! PUT VECX INTO FINAL OUTPUT
6        c(j,i)=vecx(i)

    4 continue

    cnorm=c(1,1)
    if (cnorm.eq.0.) cnorm=1.

c
c FINALLY, FLIP RIGHT SIDE TO THE LEFT OF THE LEFT SIDE:
c
    do 10 j=1,n1/2
        do 11 i=1,n2/2
            b(j,i     )=c(j,i+n2/2)/cnorm
11        b(j,i+n2/2)=c(j,i     )/cnorm
10    continue

    return
end

```